# The Generalized Hilbert Boundary Value Problem 

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## Abstract

In this work we study the Noether and the solvability theories of the generalized Hilbert boundary value problem. Our main goal is to obtain the defect numbers of that problem with a direct or an inverse linear fractional Carleman shift of order $2(\alpha(\alpha(t)) \equiv t)$ on the unit circle $\pi$. To this end we start by reducing the mentioned problem to a singular integral operator with shift. Afterwords we use the well known fact that this operator can be reduced to a singular integral operator (without shift) whose coefficients are matrix functions. Finally we compute the partial indices of such matrix functions (which, in this case, can be represented as a product of an Hermitean matrix function with negative determinant by diagonal rational matrix functions) and use these results to obtain the defect numbers of the initial problem.

In what concerns the Noether theory of the generalized Hilbert boundary value problem, we obtain results in the cases of a non-Carleman shift or of a Carleman shift of arbitrary orde
We also give examples of this problem, with the direct Carleman shift $\alpha(t)=-t$ and with the inverse Carleman shift $\alpha(t)=\frac{1}{t}$, for which the Gakhov-Coburn formulas are not valid.
Key Words: Boundary value problems; Singular integral operators; Shift; Factorization; Noether theory; Solvability theory

## Introduction

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Let }\Gamma\mathrm{ be a simple closed Lyapunov curve dividing the complex plane C i two parts D+ and D D and (
Hibert boundary value problem consists in finding a function
    \Phi+}(z)=u(x,y)+iv(x,y),z=x+iy
analytic in the domain D}\mathrm{ , whose limit values of its real and imaginary parts belong to }\mp@subsup{H}{\mu}{}(\Gamma)\mathrm{ and
satisfy on \Gamma the condition
    a(t)u(t)+b(t)uu(\alpha(t))+c(t)v(t)+d(t)v(\alpha(t))=h(t),
    Re{\mathcal{A}(t)\mp@subsup{\Phi}{}{+}(t)+\mathcal{B}(t)\mp@subsup{\Phi}{}{+}(\alpha(t))}=h(t)
with \(\mathcal{A}(t)=a(t)-i c(t)\) and \(\mathcal{B}(t)=b(t)-i d(t)\), where \(a, b, c, d, h \in H_{\mu}(\Gamma)\) are real functions. This problem is a generalization of the (classical) Hilbert boundary value problem (obtained if \(b(t)=d(t) \equiv 0)\) and was proposed by E. G. Khasabov and G. S. Litvinchuk [1]. In their papers
\([11\) and \([2]\) the Noetherity conditions and the index formula of problem (1) with a direct or an inverse Carleman shift of order \(2(\alpha(\alpha(t)) \equiv t)\) on \(\Gamma\) were obtained. In what concerns the solvability theory it was only established for some particular cases of degeneracy of boundary condition (1) which are reduced to the well known binomial boundary value problems of Hilbert and of Carleman type. Preliminaries
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$\qquad$
(i) $\operatorname{im} A=\overline{\mathrm{imA}}$ (i.e. $A$ is a normally solvable).
ii) The numbers $\alpha(A)=\operatorname{dim} \operatorname{ker} A$ and $\beta(A)=\operatorname{dim} \operatorname{coker} A\left(\operatorname{coker} A=X_{2} / \mathrm{im} A\right)$ are finite.
The subspaces ker $A$ and coker $A$ are called the defect subspaces of operator $A$ and the numbers $\alpha(A)$
and $\beta(A)$ are called the defect numbers of this operator:
Definition 2 The integer $\mathcal{I}(A)=\operatorname{ind} A=\alpha(A)-\beta(A)$ is called the index of the Noetherian operator
We also notice that an equation of the form $A x=y$ is solvable if and only if $\rho=\beta(A)$ solvability
conditions are fulfililed and, in that case it has $l=\alpha(A)$ linearly independent solutions. The number
$I=l-\rho$ is also called the index of the ( N o
satisfy the Gakhov-Coburn formulas when
$l=\max (0, I)$ and $\rho=\max (0,-I)$

## The Noether theory of an operator consists in finding a Noetherity criterion for that operator and cal

 culating its indexThe solvability theory of an operator includes the calculation of its defect numbers, the construction of bases for the defect subspaces, the problems of spectral theory, the determination of exact or approximate solutions for the corresponding equations and boundary value problems.

Definition 3 Let $\Gamma$ be a simple closed smooth contour dividing the closed complex plane into the do mains $D^{+}(\ni 0)$ and $D^{-}(\ni \infty)$. A factorization of a non-singular matrix function $G(t)$ relative to the contour $\Gamma$ is a representation of $G$ in the form

$$
G(t)=G^{+}(t) \Lambda(t) G^{-}(t)
$$

where $G^{ \pm}(t)$ are the boundary values of matrix functions $G^{ \pm}(z)$, analytic and non-singular in $D^{ \pm}$
 integers wich we call
the total index of $G$.

Let $\Gamma$ be a simple oriented curve, a homeomorphism $\alpha: \Gamma \rightarrow \Gamma$ is called a shift. A homeomorphism $\alpha$ preserving (changing) the orientation on $\Gamma$ is called a direct (inverse) shift and sometimes is denoted Definition 4 A point $\tau \in \Gamma$ is called a periodic point of the shift $\alpha$ with multiplicity $k \geq 1(k \in \mathbb{N})$,
if $\alpha_{k}(\tau)=\tau$ and for $\left.k>1\right) \alpha_{i}(\tau) \neq \tau$ for $i \in\{1,2, \ldots, k-1\}$ where $\alpha_{i}(t)=\alpha\left(\alpha_{i-1}(t)\right)$ and we
agree that $(t) \equiv t$ agree that $\alpha_{0}(t) \equiv t$.
A periodic point with multiplicity one ( $k=1$ ) is called a fixed point
We denote $\mathcal{M}(\alpha, k)$ the set of periodic points of the shift $\alpha$ with multiplicity $k$,
Definition 5 A shift $\alpha$ satisfying the condition

$$
\alpha_{k}(t)=t \quad \forall t \in \Gamma,
$$

for some $k \geq 2$, is called a Carleman shift. A shift $\alpha$ satisfying the condition $\mathcal{M}(\alpha, k) \neq \Gamma$, for all
$k \in \mathbb{N}$ is called a non-Carleman shift.
The least value of $k$ for which the Carleman condition (4) is fulfilled is called the order of the shift $\alpha$. Considering the particular case $\Gamma=\mathbb{T}=\{t$
shift of order $2(\alpha(\alpha(t)) \equiv t)$ of the form

$$
\begin{equation*}
\alpha(t)=\frac{t-\beta}{\bar{\beta} t-1}, \quad t \in \mathbb{T}, \quad \beta \in \mathbb{C} \backslash \pi \tag{5}
\end{equation*}
$$

which preserves or changes the orientation depending on $|\beta|<1$ or $|\beta|>1$, respectively. The shift (5) admits the factorization
$\alpha(t)=\alpha^{+}(t) t^{\mu} \alpha^{-}(t)$, (6) where $\alpha^{+}(t)=\lambda(\bar{\beta} t-1)^{-1}, \alpha^{-}(t)=\lambda^{-1}(t-\beta) t^{-1}, \lambda=\sqrt{1-|\beta|^{2}}, \mu=1$, if $|\beta|<1$ and $\alpha^{+}(t)=(i \lambda)^{-1}(t-\beta), \alpha^{-}(t)=i \lambda t(\bar{\beta} t-1)^{-1}, \lambda=\sqrt{|\beta|^{2}-1}, \mu=-1$, if $|\beta|>$
Noether theory of the generalized Hilbert boundary value problem

In [1] and [2], E. G. Khasabov and G. S. Litvinchuk used an integral representation to obtain the
 $2(\alpha(\alpha(t)) \equiv t)$ on $\Gamma$ (see also [5]). Here we use the function $\nu_{\alpha}$ defined in section $1.3-1.6$ of $\S 1$ of
chapter 2 of $[4]$ and the results exposed there to obtain the Noether theory of problem (2) when $\alpha$ is a chan-Carleman shift or a Carleman shift of any order $k \geq 2$.
We start by considering the case of a direct shift $\alpha$ with a non-empty arbitrary set of periodic points with multiplicity $k \geq 1$ (i.e. for some $k \geq 1, \mathcal{M}(\alpha, k) \neq \emptyset)$

Theorem 6 In the above conditions, the generalized Hilbert boundary value problem is Noetherian if $\nu_{\alpha}(\mathcal{A}, \mathcal{B}) \neq 0$ on $\Gamma$
in that case, the inder formula for this problem is

We note that the case of a direct Carleman shift of any order $k \geq 2$ is included in the previous theorem.
Is remains to consider the case of an inverse shift. In the general case of an inverse shift $\alpha$ such that $\alpha_{2}$ the curve $\Gamma$, the Noether theory of the generalized Hilbert problem is given by the following
Theorem 7 For a shift in the above conditions, the generalized Hilbert boundary value problem is Noetherian if
$\nu_{\alpha_{2}}(\mathcal{A}(t) \overline{\mathcal{A}(\alpha(t)),}, \mathcal{B}(\alpha(t)) \overline{\mathcal{B}(t)}) \neq 0$ sobre $\Gamma$,
and, if this condition holds, the index of problem (2) is given by the formula
$I=-\operatorname{Ind}_{\Gamma} \nu_{\alpha_{2}}(\mathcal{A}(t) \overline{\mathcal{A}(\alpha(t))}, \mathcal{B}(\alpha(t)) \overline{\mathcal{B}(t)})+1$
Solvability theory of the generalized Hilbert boundary value problem

In this subsection we obtain the defect numbers $l$ and $\rho$ of the generalized Hilbert boundary value Carleman shift of order $2(\alpha(\alpha(t)) \equiv t)$ of the form (5). We start by introducing some identities which will be used in this section. Put
$\begin{aligned} & \Delta(t)=\mathcal{A}(t) \mathcal{A}(\alpha(t))-\mathcal{B}(t) \mathcal{B}(\alpha(t)) \\ & \theta(t)=\frac{\mathcal{A}(t) \mathcal{A}(\alpha(t))}{\mathcal{B}}, \\ &\mathcal{B}(t) \mathcal{B}(t))\end{aligned}$
$\begin{aligned} & \theta(t)=\mathcal{A}(t(\mathcal{A}(\alpha(t)) \\ & V(t)=\mathcal{B}(t) \mathcal{\mathcal { A }}(\alpha(t)) \\ & \mathcal{B}(t) \mathcal{H}(\alpha(t)),\end{aligned}$
It can be directly verified that $\operatorname{Re} V=0$, consequently $V=i V_{0} \operatorname{com} V_{0}$ função real ( $V_{0}=-i V$ ).
We observe that, using the notations introduced above, the Noetherity condition of problem (2) takes
The case of a direct linear fractional Carleman shift
In this subsection $\alpha$ is a direct linear fractional Carleman shift of order $2(\alpha(\alpha(t)) \equiv t)$ of the form
$(5)($ with $|\beta|<1)$, and we introduce the weighted shift operator $(U \varphi)(t)=u(t) \varphi(\alpha(t))$, where $u(t)=-\alpha^{+}(t)=-\lambda(\bar{\beta} t-1)^{-1}($ see $(6))$.
We start by observing that, if $\Gamma=\pi$, the study of the solvability theory of problem (2) can be reduced to the study of the solvability theory of the singular integral operator with shift

$$
T=\left(\mathcal{A}(t) I+u^{-1}(t) \mathcal{B}(t) U\right) P_{+}-\left(\bar{t} \overline{\mathcal{A}(t)} I+\alpha(t) u^{-1}(t) \overline{\mathcal{B}(t)} U\right) P_{-}
$$

with $P_{ \pm}=\frac{1}{2}(I \pm S)$, where $S$ is the operator of singular integration and $I$ is the identity operator. On the other hand, as $U S=S U$, the solvability theory of operator $T$ can be obtained from the solvability theory of the singular integral operator without shift

$$
\begin{gathered}
M=A P_{+}+B P_{-}, \quad \text { with } \\
A=\left(\begin{array}{cc}
\mathcal{A}(t) & u^{-1}(t) \mathcal{B}(t) \\
u(t) \mathcal{B}(\alpha(t)) & \mathcal{A}(\alpha(t))
\end{array}\right) \text { and } B=\left(\begin{array}{cc}
-t \overline{\mathcal{A}(t)} & -\alpha(t) u^{-1}(t) \overline{\mathcal{B}(t)} \\
-t u(t) \overline{\mathcal{B}(\alpha(t))} & -\alpha(t) \overline{\mathcal{A}(\alpha(t))}
\end{array}\right)
\end{gathered}
$$

Hence, it is enough to study the solvability theory of operato

$$
\widetilde{M}=P_{+}+C P_{-},
$$

with
$C=A^{-1} B=\frac{1}{\Delta(t)}\left(\begin{array}{cc}i & 0 \\ 0 & -i\end{array}\right)\left(\begin{array}{cc}u^{-1}(t) & 0 \\ 0 & 1\end{array}\right) C_{0}\left(\begin{array}{cc}1 & 0 \\ 0 & u(t)\end{array}\right)\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)\left(\begin{array}{cc}t & 0 \\ 0 & \alpha(t)\end{array}\right), \quad(8)$ where $C_{0}=\left(\begin{array}{cc}\frac{V_{0}(t)}{i \theta(t)} & i \theta(t) \\ -V_{0}(\alpha(t))\end{array}\right)$ is an Hermitean matrix function with negative determinant. Then by virtue of Theorem 2 of $[6]$ the partial indices of the matrix function $C_{0}$ are $m(\geq 0)$ and $-m$. Now, using $(8)$, and taking into account the factorization $\alpha(t)=\alpha^{+}(t) t \alpha^{-}(t)($ see $(6))$ and the idenlity $u(t)=-\alpha^{+}(t)$, we obtain that the partial indices of the matrix function $C$ are
where

$$
\kappa=\frac{1}{2 \pi}\{\arg \overline{\Delta(t)}\}_{\Gamma}=\operatorname{Ind} \overline{\Delta(t)} .
$$

Finally, using Theorem 2.10 of $[3]$, we obtain the following result about the solvability theory of the generalized Hilbert boundary value problem (2) with a direct linear fractional Carleman shift on $\pi$.

Theorem 8 In the above conditions, the numbers $l$, of linearly independent solutions, and $\rho$, of solvbility conditions of the generalized Hibert bowdary wine problem (2) with a direct linear fractional Carleman shift of oder 2 on $T$ are given by

1. $l=\max (0, \kappa+1) \quad$ and $\quad \rho=\max (0,-\kappa-1)$, if $\kappa_{1} \leq 0$ or $\kappa_{2}>0$,
2. $l=\operatorname{dim} \operatorname{ker} T=\frac{\kappa_{1}}{2}=\frac{\kappa+m+1}{2}$ and $\rho=\frac{-\kappa+m-1}{2}$, if $\kappa_{1}>0, \kappa_{2} \leq 0$ and $\kappa_{1}$ is even,
3. $l=\left[\frac{\kappa}{2}+\frac{1-\varepsilon}{4}\right]=\left\{\begin{array}{ll}\frac{k+m}{\varepsilon}, & \text { se } \varepsilon=1 \\ \frac{k+m}{2}+1, & \text { se } \varepsilon=-1\end{array}\right.$ and $\rho=\left\{\begin{array}{ll}\frac{-k+m}{k+m}-1 & , \text { se } \varepsilon=1 \\ \frac{k,}{2} & \text { se } \varepsilon=-1\end{array}\right.$, if $\kappa_{1}>0, k_{2} \leq 0$

Now we present an example (see also [2], and [5], example 23.1) of a generalized Hilbert boundary value problem for which the Gakhov-Coburn formulas (3) are not always valid

Example 9 Let us consider the generalized Hilbert problem
$\operatorname{Re}\left\{\cos (k \varphi) \Phi^{+}(t)-i \sin (k \varphi) \Phi^{+}(-t)\right\}=h(t) \quad, \quad t=e^{i \varphi} \in \mathbb{\pi}$, with $k \in \mathbb{Z} . \quad$ (9)
In this case (see (2))
$\alpha(t)=-t, \mathcal{A}(t)=\cos (k \varphi)=\operatorname{Ret}^{-k}, \mathcal{B}(t)=-i \sin (k \varphi)=i \operatorname{IIm} t^{-k}$
$\Delta(t)=(-1)^{k}, \theta(t)=(-1)^{k} \frac{t^{2 k}+t^{-2 k}}{2}, V(t)=-(-1)^{k} \frac{t^{2 k}-t^{-2 k}}{2}$
From this we conclude that problem $(9)$ is Noetherian $(\Delta(t) \neq 0)$ and, as $\kappa=\operatorname{lnd} \pi \Delta(t)=0$, its index is given by $I=\kappa+1=1$. Besides, for this example, according to $(8)$
circulant matrix exposed in $[7]$, p. 159 , the matrix function $C$ comes

$$
C=-\frac{1}{2}\left(\begin{array}{cc}
1 & 1 \\
1 & -1
\end{array}\right)\left(\begin{array}{cc}
t^{2 k+1} & 0 \\
0 & t^{-2 k+1}
\end{array}\right)\left(\begin{array}{cc}
1 & -1 \\
1 & 1
\end{array}\right) .
$$

Hence, the partial indices of the matrix function $C$ are given by
$k_{1}=2|k|+1 \quad$ e $\quad k_{2}=-2|k|+1$,
Finally, using Proposition 2.2 of [3] to compute the numb
$k<0, \varepsilon=1$. Thus, by virtue of Theorem 8 , we obtain

$$
\left\{\begin{array}{lll}
l=k+1 \text { and } \rho=k, & \text { if } k \geq 0 \\
l=-k & \text { and } \rho=-k-1, & \text { if } k<0
\end{array}\right.
$$

The case of an inverse linear fractional Carleman shift h what follows $\alpha$ is an inverse linear fracional Carleman shift $\alpha$ of order $2(\alpha(\alpha(t)) \equiv t)$ of the form 5) ( with $|3|>1$ ).

Now we consider the operator $U$ defined by $(U \varphi)(t)=u(t) \varphi(\alpha(t))$, with $u(t)=\alpha^{-}(t) t^{-1}=\frac{i \lambda}{\beta t-1}$ see (6)).
As we have seen in the previous subsection, the study of the solvability theory of problem (2) can be educed to the study of the solvability theory of operator ( 7 ).
Analogously to what happened in the case of a direct shift, now, as $U S=-S U$, the solvability theory operator ( 7 ) can be obtained from the solvability theory of the singular integral operator without $\widetilde{M}=P_{+}+C P_{-}$,
with

$$
C=A^{-1} B=\frac{1}{\alpha(t) \overline{\theta(t)}}\left(\begin{array}{cc}
i \alpha(t) u^{-1}(t) & 0  \tag{10}\\
0 & -i
\end{array}\right) C_{0}\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right)\left(\begin{array}{cc}
t u(t) & 0 \\
0 & 1
\end{array}\right)
$$

where, and $C_{0}=\left(\begin{array}{cc}V_{0}(t) & \overline{i \Delta(t)} \\ i \Delta(t) & V_{0}(\alpha(t))\end{array}\right)$
ow, by virtue of Theorem 2 of $[6]$ we have that the partial indices of the matrix function $C_{0}$ are $m(\geq 0)$ and $-m$.
Using the factorization $\alpha(t)=\alpha^{+}(t) t^{-1} \alpha^{-}(t)$ and the fact that $u(t)=\alpha^{-}(t) t^{-1}$, from (10) we conclude that the partial indices of the matrix function $C$ are
where

## $\kappa=\frac{1}{2 \pi}\{\arg \theta(t)\}_{\Gamma}=\operatorname{Ind} \pi \theta(t)$

Finally, by virtue of Theorem 1 of section 21 of [5] (p. 239), we obtain the following result on the defect numbers of the generalized Hilbert boundary value problem (2).
heorem 10 In the above conditions, puting $I=\operatorname{ind} T=\kappa+1$, the numberst, of inearly independen an inverse linear fractional Carleman shift of order 2 of the form (5) on $\mathbb{T}$ are given by:

## $l=\max (0, I) \quad$ and $\quad \rho=\max (0,-I), i f \kappa_{1} \leq 0$ or $\kappa_{2}>0$,

2. $l=\frac{\kappa_{1}}{2}=\frac{\kappa+1+m}{2}$ and $\rho=\frac{-k-1+m}{2}$, if $\kappa_{1}>0$ and $k_{2} \leq 0$ and $\kappa_{1}$ is even.
3. $l=\frac{\kappa_{1}-\varepsilon}{2}=\frac{\kappa+1+m-\varepsilon}{2}$ and $\rho=\frac{-k-1+m-\varepsilon}{2}$, if $\kappa_{1}>0$ and $\kappa_{2} \leq 0$ and $\kappa_{1}$ is odd, where $\varepsilon= \pm 1$.

The previous theorem is also valid if we consider the shift $\alpha(t)=\frac{1}{t}$. First of all we observe that this is an inverse Carleman shift $(\alpha(\alpha(t)) \equiv t)$ on $\mathbb{T}$, but it is not a particular case of a linear fractional Carleman shift of the form $(5)$. Afterwords, if together with this shift we consider the so called $f i p$
operator $U$, defined by $(U \varphi)(t)=\frac{1}{1} \varphi\left(\frac{1}{)}\right.$ it is easy to show that all the results obtained above for the case of an inverse linear fractional Carleman shift of the form (5) remain valid in these conditions. In fact, it is enough to replace, in the above procedure, $u(t)$ for $u(t)=t^{-1}$ and to observe that in this ase, as $\alpha(t)=t^{-1}$, we have that $\alpha^{+}(t)=\alpha^{-}(t)=1($ see $(6))$.
Finally we present an example (see also [2] and [5] (example 23.2)) of the generalized Hilbert boundfulfill the Gakhov-Coburn formulas (3).

Example 11 We consider the generalized Hilbert boundary value problen

$$
\begin{equation*}
\operatorname{Re}\left\{t^{-k}\left[u(t)+i v\left(\frac{1}{t}\right)\right]\right\}=h(t) \quad, \quad t=e^{i \varphi} \in \mathbb{T}, \text { with } k \in \mathbb{Z} \tag{11}
\end{equation*}
$$

Now (see (2))

$$
\begin{aligned}
& \alpha(t)=\frac{1}{t}, \mathcal{A}(t)=\cos (k \varphi)=\operatorname{Ret}^{-k}, \mathcal{B}(t)=-i \sin (k \varphi)=i \operatorname{Im} t^{-k} \\
& \theta(t)=1, \Delta(t)=\frac{t^{2 k}+t^{-2 k}}{2}, V(t)=-\frac{t^{2 k}-t^{-2 k}}{2} .
\end{aligned}
$$

particular we observe that, as $\theta(t)=1 \neq 0$, problem (11) is Noetherian. Besides, $\kappa=\operatorname{Ind} \tau \theta(t)=0$ this case, from (10) and using the factorization of a circulant matrix (see e.g. [7], p. 159), we obtain

$$
C=\frac{1}{2}\left(\begin{array}{cc}
-1 & -1 \\
1 & -1
\end{array}\right)\left(\begin{array}{cc}
t^{2 k+1} & 0 \\
0 & t^{-2 k+1}
\end{array}\right)\left(\begin{array}{cc}
-1 & 1 \\
-1 & -1
\end{array}\right)
$$

$\kappa_{1}=2|k|+1 \quad$ e $k_{2}=-2|k|+1$,
Iremains to obtain the number $\varepsilon$. Proceeding as in example 9 we obtain that, if $k \geq 0, \varepsilon=-1$ and if $<0, \varepsilon=1$. Thus, applying Theorem 10 , we conclude that
i) $l=\kappa+1$ and $\rho=\kappa$, if $k \geq 0$,
(ii) $l=-k$ and $\rho=-k-1$, if $k<0$.

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